

# Chapter 30

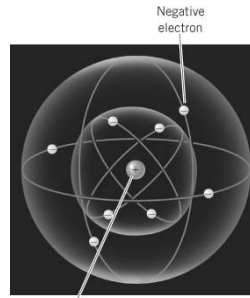
## Atom Physics



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### 30.1 Rutherford Scattering and the Nuclear Atom



Model of an atom-the recent model

Nuclear radius –  $r \approx 10^{-15}$  m

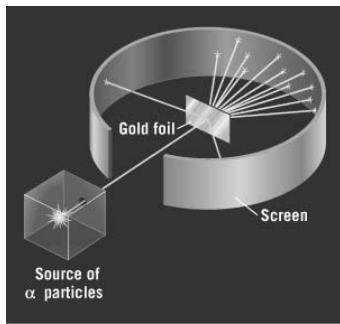
Electron's position radius  
 $r \approx 10^{-10}$  m = 1 Å

Positive nucleus  
Negative electron  
In its natural state, an atom is electrically neutral.

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### 30.1 Rutherford Scattering and the Nuclear Atom: Planetary model



Positively charged alpha particles from radioactive source bombard a thin gold foil being reflected. Hence Thomson's plum-pudding model of atom has to be rejected since only concentrated positive charges can cause deflection due to  $F_E$

Electrons must move around nucleus, else it will be electrically attracted to the positive nucleus. But accelerating electrons radiate and hence lose energy via EM wave radiation. So???

A Rutherford scattering experiment.

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### 30.1 Rutherford Scattering and the Nuclear Atom

#### Conceptual Example 1 Atoms are Mostly Empty Space

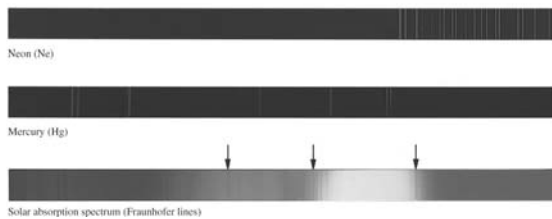
In the planetary model of the atom, the nucleus (radius =  $10^{-15}$ m) is analogous to the sun (radius =  $7 \times 10^8$ m). Electrons orbit (radius =  $10^{-10}$ m) the nucleus like the earth orbits (radius =  $1.5 \times 10^{11}$ m) the sun. If the dimensions of the solar system had the same proportions as those of the atom, would the earth be closer to or farther away from the sun than it actually is?

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### 30.2 Line Spectra

The individual wavelengths emitted by two gases and the continuous spectrum of the sun.



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### 30.2 Line Spectra-bright fringes emitted by excited gases

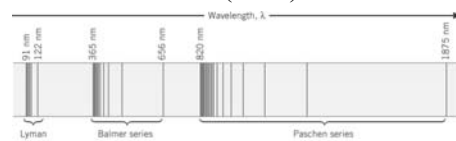
#### The Line Spectrum of Hydrogen

$R$  is Rydberg constant  
 $R = 1.097 \times 10^7 \text{ m}^{-1}$

Lyman series: transition to ground state  $\frac{1}{\lambda} = R \left( \frac{1}{1^2} - \frac{1}{n^2} \right)$   $n = 2, 3, 4, \dots$

Balmer series: transition to 1<sup>st</sup> excited state  $\frac{1}{\lambda} = R \left( \frac{1}{2^2} - \frac{1}{n^2} \right)$   $n = 3, 4, 5, \dots$

Paschen series: transition to 2<sup>nd</sup> excited state  $\frac{1}{\lambda} = R \left( \frac{1}{3^2} - \frac{1}{n^2} \right)$   $n = 4, 5, 6, \dots$



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**30.3 The Bohr Model of the Hydrogen Atom-explaining the observed experimental line spectrum**

In the Bohr model, a photon is emitted when the electron drops from a larger, higher-energy orbit to a smaller, lower energy orbit.

In a hydrogen atom, electrons' total (kinetic & potential) energy can have only certain allowed values corresponding different orbits (stationary state or orbits) of the electrons around the nucleus & that electrons do not radiate while orbiting

Photons are only emitted when electrons changes orbits from larges to smaller orbits. The energy of the photon is

$$\Delta E = E_i - E_f = hf = \frac{hc}{\lambda}$$

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**30.3 The Bohr Model of the Hydrogen Atom**

THE ENERGIES AND RADII OF THE BOHR ORBITS

Assume the electrons are not radiating while orbiting the nucleus. The centripetal force is provided by the electrical force

$$F_{12} = \frac{kq_1q_2}{r^2}$$

$$E = KE + EPE$$

$$= \frac{1}{2}mv^2 - \frac{kZe^2}{r}$$

$$= -\frac{kZe^2}{2r}$$

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**30.3 The Bohr Model of the Hydrogen Atom**

ENERGY LEVEL DIAGRAMS

Total energy,  $E$

$n = \infty$ , Electron removed from atom

$n = 5$

$n = 4$

$n = 3$  Excited states

$n = 2$

$n = 1$ , Ground state

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**30.3 The Bohr Model of the Hydrogen Atom**

THE LINE SPECTRA OF THE HYDROGEN ATOM

Total energy,  $E$

$n = \infty$

$n = 5$

$n = 4$

$n = 3$

$n = 2$  Balmer series

$n = 1$  Lyman series

$$\frac{1}{\lambda} = \frac{2\pi^2mk^2e^4}{h^3c} (Z^2) \left( \frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

$n_i, n_f = 1, 2, 3, \dots \quad n_i > n_f$

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**30.5 The Quantum Mechanical Picture of the Hydrogen Atom**

**Example 5 The Bohr Model Versus Quantum Mechanics**

Determine the number of possible states for the hydrogen atom when the principal quantum number is (a)  $n=1$  and (b)  $n=2$ .

$n = 1 : \ell = 0 : m_\ell = 0$

$m_s = +\frac{1}{2}$

$m_s = -\frac{1}{2}$

State			
$n$	$\ell$	$m_\ell$	$m_s$
1	0	0	$+\frac{1}{2}$
1	0	0	$-\frac{1}{2}$

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**30.5 The Quantum Mechanical Picture of the Hydrogen Atom**

$n = 2$

$\ell = 1$

$m_\ell = +1$

$m_s = +\frac{1}{2}$

$m_s = -\frac{1}{2}$

$m_\ell = 0$

$m_s = +\frac{1}{2}$

$m_s = -\frac{1}{2}$

$m_\ell = -1$

$m_s = +\frac{1}{2}$

$m_s = -\frac{1}{2}$

$\ell = 0$

$m_\ell = 0$

$m_s = +\frac{1}{2}$

$m_s = -\frac{1}{2}$

State			
$n$	$\ell$	$m_\ell$	$m_s$
2	1	+1	$+\frac{1}{2}$
2	1	+1	$-\frac{1}{2}$
2	1	0	$+\frac{1}{2}$
2	1	0	$-\frac{1}{2}$
2	1	-1	$+\frac{1}{2}$
2	1	-1	$-\frac{1}{2}$
2	0	0	$+\frac{1}{2}$
2	0	0	$-\frac{1}{2}$

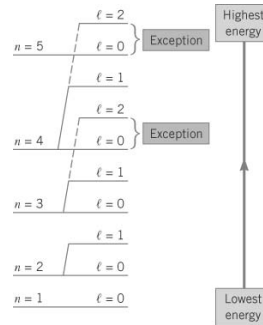
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30.5 The Quantum Mechanical Picture of the Hydrogen Atom

**Conceptual Example 6 The Bohr Model Versus Quantum Mechanics**

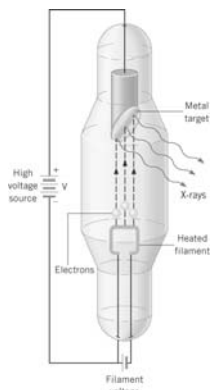
Consider two hydrogen atoms. There are no external magnetic fields present, and the electron in each atom has the same energy. According to the Bohr model and to quantum mechanics, is it possible for the electrons in these atoms (a) to have zero orbital angular momentum and (b) to have different angular momenta?

30.6 The Pauli Exclusion Principle and the Periodic Table of the Elements



Generally, the energy increases with increasing  $n$ . There are exceptions to the general rule.

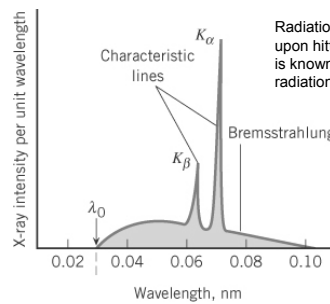
30.7 X-Rays



Electrons are emitted from a heated filament and accelerated through a large voltage.

When they strike the target (platinum or molybdenum, X-rays (photons when electrons undergo transition from higher energy states to lower energy states) are emitted.

30.7 X-Rays



Radiation produced due to deceleration upon hitting the target is a continuum and is known as Bremsstrahlung (braking radiation)

The sharp peaks are called **characteristic X-rays (X-ray photons)** because they are characteristic of the target material. The  $K_\alpha$  (from  $n=2$  to  $n=1$ ) and  $K_\beta$  ( $n=3$  to  $n=1$ ) lines involve transition of electrons from an excited state filling in the vacancy created in the K shell ( $n=1$ ) after the electron from that shell is ejected by the incoming electrons.

$\lambda_0$  is the cutoff wavelength which depends ONLY on the energy of the impinging electrons.

$$KE = eV = hf = \frac{hc}{\lambda_0} \quad \text{So, } \lambda_0 = \frac{hc}{eV}$$

30.7 X-Rays

